

1. Consider the system in polar coordinates. Show that there are numbers  $0 < r_1 < r_2$  so that  $r_1 \leq r \leq r_2$  is a trapping region. Show that the system has a nontrivial periodic trajectory.

$$
\dot{r} = r(3 - 2r^2 + r^2 \sin^2 \theta)
$$

$$
\dot{\theta} = 1
$$

Using the fact that  $0 \le \sin^2 \theta \le 1$  we find that

$$
r(3 - 2r^{2}) \le r = r(3 - 2r^{2} + r^{2}\sin^{2}\theta) \le r(3 - r^{2})
$$

It follows that if one takes any  $0 < r_1 < \sqrt{\frac{3}{2}}$  then  $\dot{r} > 0$  whenever  $r = r_1$  and flow is outward through the circle  $r = r_1$ . If one takes any  $\sqrt{3} < r_2$  then  $\dot{r} < 0$  whenever  $r = r_2$  so flow is inward through the circle  $r = r_2$ . Thus the annulus  $r_1 \le r \le r_2$  is a trapping region. Also, observe that  $\theta \neq 0$  for all  $r > 0$  so that there are no fixed points in the annulus. Since the vector field is smooth, we may apply the Poincaré-Bendixson Theorem, which asserts that any trajectory starting in the trapping region tends to a nonconstant limit cycle, which is a nontrivial periodic trajectory in the trapping annulus.

2. Consider the system in polar coordinates. The system undergoes a bifurcation as the parameter  $\mu > 0$  passes the critical value  $\mu_c$ . Find the value and sketch the phase portraits for  $\mu < \mu_c$ ,  $\mu = \mu_c$  and  $\mu > \mu_c$ . What kind of bifurcation is this?

$$
\dot{r} = r(1 - r^2)
$$

$$
\dot{\theta} = 1 + \mu \cos \theta
$$

The zeros of the  $\dot{r} = r(1 - r^2)$  occur at  $r = 0$  and  $r = 1$ . The origin is an unstable rest point and the circle  $r = 1$  is an invariant set. Since  $\dot{r} > 0$  in  $0 < r < 1$  and  $\dot{r} < 0$  if  $1 < r$  we see that the  $r = 1$  is an attractive rest point. Flows starting away from the origin or unit circle tends toward the unit circle.  $\theta = 1 + \mu \cos \theta \ge 1 - \mu$  for all  $\theta$  so  $\theta(t)$  is strictly increasing for  $0 < \mu < 1$  and, except for the origin, the flow approaches the limit cycle  $r = 1$ . At  $\mu = 1$ , an infinite time bifurcation occurs:  $\theta(t)$  stops increasing at one point where  $1 + \cos \theta = 0$ or  $\theta = \pi$ . For  $\mu > 1$  there are two rest points  $\theta = \phi_{\pm}$  with  $0 < \phi_{-} < \pi < \phi_{+}$  which solve

 $0 = 1 + \mu \cos \theta$ , namely,  $\phi_{\pm} = \arccos(1/\mu)$ .

Then  $\dot{\theta} < 0$  for  $\phi_- < \theta < \phi_+$  and  $\dot{\theta} > 0$  otherwise. Radially flow approaches the  $r = 1$ circle but  $\theta$  < 0 for initial angles  $\phi$ <sub>-</sub> <  $\theta$ <sub>0</sub> <  $\phi$ <sub>+</sub> and positive otherwise. This makes  $(r, \theta) = (1, \phi_{-})$  a stable node and  $(1, \phi_{+})$  an unstable saddle.



Figure 1: Plots with  $\mu=.5,1,1.3$  using 3D-XPlorMath<br>@

- 3. Answer the following questions about periodic trajectories.
	- (a) Show that this equation has no nontrivial periodic solutions.

$$
\ddot{x} + x^2 \dot{x} + x = 0.
$$

Viewing this as a spring equation with nonnegative drag depending on velocity, one expects that energy decay under the flow. Written as a system

$$
\begin{aligned}\n\dot{x} &= y\\ \n\dot{y} &= -x - xy^2\n\end{aligned}
$$

Computing

$$
\frac{d}{dt}E = \frac{d}{dt}\frac{x^2 + y^2}{2} = x\,\dot{x} + y\,\dot{y} = xy + y(-x - xy^2) = -x^2y^2 \le 0.
$$

If  $(x, y) \neq (0, 0)$ , the flow does not stop at  $x = 0$  because  $\dot{x} = y \neq 0$  nor at  $y = 0$ because then  $\dot{y} = -x \neq 0$ . Otherwise  $E < 0$  so the energy is strictly decreasing function of time. It follows that there cannot be nontrivial periodic trajectories because the energy cannot be periodic.

(b) Show that this equation has a nontrivial periodic solution.

$$
\ddot{x} + x(\dot{x})^2 + x = 0.
$$

Written as a system

$$
\begin{aligned}\n\dot{x} &= y &= f(x, y) \\
\dot{y} &= -x - x^2 y &= g(x, y)\n\end{aligned}
$$

This system is reversible because  $f(x, -y) = -f(x, y)$  is odd in y and  $g(x, -y) =$  $g(x, y)$  is even in y. Its only zero is when  $0 = f(x, y) = y$  hence  $y = 0$  and when  $0 = g(x, 0) = -x$  so  $x = 0$  as well. Linearizing at zero we find

$$
J(0,0) = \begin{pmatrix} 0 & 1 \\ -1 - 2xy & -x^2 \end{pmatrix} \Big|_{(x,y)=(0,0)} = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}
$$

which is a matrix whose eigenvalues are  $\pm i$ . The origin is a center for the linearized equation. By the theorem about centers for reversible systems, because the vector field is smooth and because the reversible system has centers for the linearization at a rest point, then the nonlinear system, too, has centers at the rest point. In particular, the origin in this problem is surrounded by closed nontrivial trajectories.

4. The system undergoes a bifurcation as the parameter  $\mu$  passes the critical value  $\mu_c$ . Find the value. Find the critical points and determine their stability. Sketch the phase portraits for  $\mu < \mu_c$ ,  $\mu = \mu_c$  and  $\mu > \mu_c$ . What kind of bifurcation is this?

$$
\begin{aligned}\n\dot{x} &= -x + y + y(\mu - y) &= f(x, y) \\
\dot{y} &= y(\mu - y) &= g(x, y)\n\end{aligned}
$$

The rest points are at solutions of  $0 = g(x, y) = y(\mu - y)$  which is at  $y = 0$  or  $y = \mu$ . If  $y = 0$  then  $0 = f(x, 0) = -x$  so  $x = 0$ . If  $y = \mu$  then  $0 = f(x, \mu) = -x + \mu$  so  $x = \mu$ . Computing the Jacobian we find

$$
J(x,y) = \begin{pmatrix} -1 & 1+\mu-2y \\ 0 & \mu-2y \end{pmatrix}; \ J(0,0) = \begin{pmatrix} -1 & 1+\mu \\ 0 & \mu \end{pmatrix}; \ J(\mu,\mu) = \begin{pmatrix} -1 & 1-\mu \\ 0 & -\mu \end{pmatrix}
$$

At  $(0,0)$ , the eigenvalues are  $-1$  and  $\mu$  with corresponding eigenvectors  $\begin{pmatrix} 1 \\ 0 \end{pmatrix}$  and  $\begin{pmatrix} 1 \\ 1 \end{pmatrix}$ . At  $(\mu, \mu)$ , the eigenvalues are  $-1$  and  $-\mu$  with corresponding eigenvectors  $\begin{pmatrix} 1 \\ 0 \end{pmatrix}$  and  $\begin{pmatrix} 1 \\ 1 \end{pmatrix}$ . Thus for  $\mu > 0$ , the rest point  $(0, 0)$  is a saddle and the rest point  $(\mu, \mu)$  is a stable node. The points swap roles if  $\mu < 0$  when the rest point  $(0,0)$  is a stable node and the rest point  $(\mu, \mu)$  is a saddle. Thus a transcritical bifurcation occurs at  $(0, 0)$  when  $\mu = \mu_c = 0$ .



Figure 2: Plots with  $\mu = .5, 0, -.5$  using 3D-XPlorMath $\odot$ 

## 5. Consider the system

$$
\begin{aligned}\n\dot{x} &= \mu x - y + (\mu + 1)x^2 - xy &= f(x, y) \\
\dot{y} &= x + x^2 &= g(x, y)\n\end{aligned}
$$

(a) The system undergoes a Hopf Bifurcation when the parameter  $\mu$  passes a critical value  $\mu_c$ . What is this critical value? What are the rest points? Can you tell if the bifurcation is subcritical or supercritical?

The rest points satisfy  $0 = g(x, y) = x(1 + x)$  so  $x = 0$  or  $x = -1$ . If  $x = 0$  then  $0 = f(0, y) = -y$  implies  $y = 0$ . If  $x = -1$  at a rest point then  $0 = f(-1, y) = 0$  $-\mu - y + (\mu + 1) + y = 1$  has no solution. Thus (0,0) is the only rest point for all  $\mu$ . Computing the Jacobian we find

$$
J(x,y) = \begin{pmatrix} \mu + 2(1+\mu)x - y & -1 - x \\ 1 + 2x & 0 \end{pmatrix}; \qquad J(0,0) = \begin{pmatrix} \mu & -1 \\ 1 & 0 \end{pmatrix}
$$

Its characteristic equation is

$$
(\mu - \lambda)(-\lambda) + 1 = \lambda^2 - \mu\lambda + 1 = 0
$$

whose solutions are

$$
\lambda = \frac{\mu \pm \sqrt{\mu^2 - 4}}{2}
$$

For  $|\mu| < 2$ , as  $\mu$  crosses  $\mu_c = 0$ , the eigenvalues are complex with real part  $\mu/2$ . Thus the Hopf bifurcation occurs at  $\mu_c = 0$  when the conjugate eigenvalues cross the imaginary axis. If  $-2 < \mu < 0$ , the origin is a stable spiral if and if  $0 < \mu < 2$ , and unstable spiral. For  $\mu < -2$ , the origin is a stable node (both eigenvalues are negative) and if  $\mu > 2$  and an unstable node (both eigenvalues are positive).

Consider the energy. For trajectories

$$
\frac{d}{dt}E = x\dot{x} + y\dot{y} = x(\mu x - y + (\mu + 1)x^2 - xy) + y(x + x^2) = \mu x^2 + (1 + \mu)x^3
$$

so at least if  $\mu = -1$  the energy is strictly decreasing (the flow doesn't stop when  $x = 0$  because then  $\dot{x} = -y$ ) hence no periodic solutions. This suggests that periodic solutions occur when  $\mu > 0$  when the origin is unstable. Thus the bifurcation is supercritical.

(b) What is the index at the origin of the vector field in part (a)? Does it depend on  $\mu$ ? Why? [Hint: Does it help you to know that a Hopf Bifurcation occurs?]

We know that a Hopf Bifurcation occurs, so that for some  $\mu > 0$  there is a periodic solution  $C$  that surrounds the origin. Since there are no rest points other than the origin, the index of  $C$  gives the index at the origin. Because the vector field is tangent to C,  $I_C = 1$ , which is the index at the origin for such  $\mu$ .

The index is the same for every  $\mu$ . We know that the origin is the only rest point for any  $\mu$  so the vector field never vanishes away form the origin, and we know that the vector field varies continuously as  $\mu$  is varied. Hence the total angle change of the vector field on a fixed loop C around the origin varies continuously. But as the angle change is an integer multiple of  $2\pi$ , it has to be constant. The index is confirmed for any  $\mu$  such that  $|\mu| > 2$  since the origin is a node or for any  $0 < |\mu| < 2$  when the origin is a spiral, both of whose index is one.